

# The Crystal Structures of Barium Hexahydroxocuprate(II) and Strontium Hexahydroxocuprate(II), $\text{Ba}_2[\text{Cu}(\text{OH})_6]$ and $\text{Sr}_2[\text{Cu}(\text{OH})_6]$

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The crystal structure of  $\text{Ba}_2[\text{Cu}(\text{OH})_6]$  has been determined by Patterson and Fourier synthesis with data collected on an automated four-circle diffractometer. The crystals are monoclinic, space group  $P2_1/c$  with  $a = 6.030$  (2),  $b = 6.440$  (2),  $c = 10.115$  (2) Å,  $\beta = 124.03$  (1)°. The structure was refined by full-matrix least-squares techniques to a conventional  $R$  index of 0.028 for 955 independent reflexions. Corrections were made for absorption, anomalous dispersion and isotropic secondary extinction. The structure contains isolated highly distorted  $[\text{Cu}(\text{OH})_6]^{4-}$  octahedra which appear to be held together primarily by barium ions. The structure may also be described as built up from close packed layers (011) consisting of  $\text{Ba}^{2+}$  and  $\text{OH}^-$  ions, the  $\text{Cu}^{2+}$  ions occupying distorted octahedral sites between every two of these layers. Thus, it represents a strongly distorted version of the well known cubic  $\text{K}_2\text{PtCl}_6$ -structure, which is built up from layers containing  $\text{K}^+$  and  $\text{Cl}^-$  with  $\text{Pt}^{4+}$  at octahedral sites in a quite similar way. Average values for the Cu–O distances are 1.965 Å for four equatorial ligands and 2.805 Å for two axial ligands. The configuration of the barium–oxygen polyhedron can best be described as a distorted di-capped trigonal prism with Ba–O distances ranging from 2.642 to 2.954 Å. Single crystals of the isomorphous compound  $\text{Sr}_2[\text{Cu}(\text{OH})_6]$  have also been prepared;  $a = 5.786$  (1),  $b = 6.154$  (1),  $c = 9.744$  (2),  $\beta = 124.15$  (1)°.

## Introduction

As part of a study on correlations between crystal structures and thermal-decomposition behaviour of hydroxy complexes we have determined the structure of barium hexahydroxocuprate(II). Several hydroxo-metallates are well known (Scholder 1962), but until now there has existed little information about the structural details of these compounds. We have grown single crystals of  $\text{Ba}_2[\text{Cu}(\text{OH})_6]$  and  $\text{Sr}_2[\text{Cu}(\text{OH})_6]$  by interdiffusion of aqueous  $\text{Ba}^{2+}$ – or  $\text{Sr}^{2+}$ – and  $\text{Cu}^{2+}$  solutions in a 50% NaOH solution.

## Crystal data and experimental

Preliminary Weissenberg and precession photographs showed the crystals to be monoclinic, space group  $P2_1/c$ . Accurate unit-cell dimensions were calculated from X-ray powder photographs registered on a Guinier-de Wolff focusing camera using  $\text{Cu K}\alpha$  radiation and high purity KCl as an internal standard. These lattice parameters were refined by the method of least squares. The calculated densities for 2 formula units per cell are in agreement with the pycnometrically measured values.

Crystal data for both compounds are listed in Table 1. An investigation of the powder X-ray patterns of  $\text{Ba}_2[\text{Cu}(\text{OH})_6]$  and  $\text{Sr}_2[\text{Cu}(\text{OH})_6]$  confirms that these salts are isostructural.

A  $\text{Ba}_2[\text{Cu}(\text{OH})_6]$  single crystal with approximate dimensions  $0.19 \times 0.10 \times 0.08$  mm was chosen for data collection on an automated four-circle Picker diffractometer using  $\text{Mo K}\alpha$  radiation (pyrolytic graphite mono-

Table 1. *Crystal data*

	$\text{Ba}_2[\text{Cu}(\text{OH})_6]$	$\text{Sr}_2[\text{Cu}(\text{OH})_6]$
Crystal system	Monoclinic	Monoclinic
Systematic absences	$h0l$ , $l = 2n + 1$	$h0l$ , $l = 2n + 1$
	$0k0$ , $k = 2n + 1$	$0k0$ , $k = 2n + 1$
Space group	$P2_1/c$	$P2_1/c$
$a$	6.030 (2) Å	5.786 (1) Å
$b$	6.440 (2)	6.154 (1)
$c$	10.115 (2)	9.744 (2)
$\beta$	124.03 (1)°	124.15 (1)°
$d_{\text{calc}}$	4.51 g cm <sup>-3</sup>	3.94 g cm <sup>-3</sup>
$d_{\text{obs}}$	4.40	3.84
$Z$	2	2
$\mu$ (for $\text{Mo K}\alpha$ )	153 cm <sup>-1</sup>	225 cm <sup>-1</sup>
Colour	Blue	Blue

chromator, 002). Intensity data were collected by a conventional moving-crystal moving-counter  $\theta$ – $2\theta$  scan over a  $2\theta$  range of 2° (plus the  $\alpha_1$ – $\alpha_2$  dispersion) with a scan speed of 0.5°/min. Background counts of 40 s were taken at each side of the peak.

Two selected standard reflexions were measured after every 10 reflexions and used for an internal scaling of the data by least-squares methods. Intensities were collected up to  $2\theta \leq 60$ °. Of the 1263 reflexions measured (including standards), 955 remained after averaging equivalent reflexions and 907 were judged to be observed according to the criterion  $I \geq 3\sigma(I)$ ,  $\sigma(I)$  being the estimated standard deviation in the net intensity based on counting statistics. The intensities were reduced to  $F_o$  in the usual way by applying Lorentz and polarization factors. Absorption corrections were made by an enlarged version (Schwarzenbach, 1972) of the program ORABS 2 (Busing, 1966) and resulted in a significantly better agreement between

intensities of measured equivalent reflexions. The calculated transmission varied from 0.25 to 0.44.

All the computations, including subsequent refinement of the structure, were performed using the program system XRAY 67 (Stewart, Kundell & Baldwin, 1970), which was enlarged and modified for use on a CDC 6400-6500 computer system by Schwarzenbach (1972).

### Structure determination and refinement

Approximate positions of the heavy atoms could be deduced from a three-dimensional Patterson map. The four barium atoms are situated in a fourfold general position 4(e) of  $P2_1/c$ , with  $z \approx \frac{1}{4}$ ; the two copper atoms are in the special position 2(c). Structure factor calculations based on refined positions of these heavy atoms yielded a conventional  $R$  ( $\sum|F_o| - |F_c|/|\sum F_o|$ ) of about 0.20. A difference electron density map showed the positions of all the oxygen atoms, which were included in the subsequent cycles of full-matrix least-squares refinement. The function minimized was  $\sum w(|F_o| - |F_c|)^2$ , where  $w$  is the weight derived from counting statistics and defined as  $w = 1/\sigma^2(F_o)$ . Structure factors were calculated with the atom scattering factors tabulated in *International Tables for X-ray Crystallography* (1962) including real and imaginary parts of anomalous dispersion corrections for barium and copper.

After several cycles of least-squares refinement in which individual isotropic temperature factors were used,  $R$  decreased to 0.059 and the weighted  $R_w\{ = [\sum w(|F_o| - |F_c|)^2 / \sum w|F_o|^2]^{1/2}\}$  to 0.079. Anisotropic temperature factors were then introduced and the isotropic secondary extinction was included as a parameter in the least-squares refinement (Larson, 1967). A final cycle of full-matrix anisotropic refinement gave  $R = 0.028$  and  $R_w = 0.029$  for 955 reflexions (including those

not observed). The resulting fractional positional coordinates are listed in Table 2 and vibrational parameters in Table 3. The corresponding observed and calculated structure factors are presented in Table 4.

### Discussion of the structure

The crystal structure of  $\text{Ba}_2[\text{Cu}(\text{OH})_6]$  contains highly distorted octahedra of  $[\text{Cu}(\text{OH})_6]^{4-}$ . These octahedra

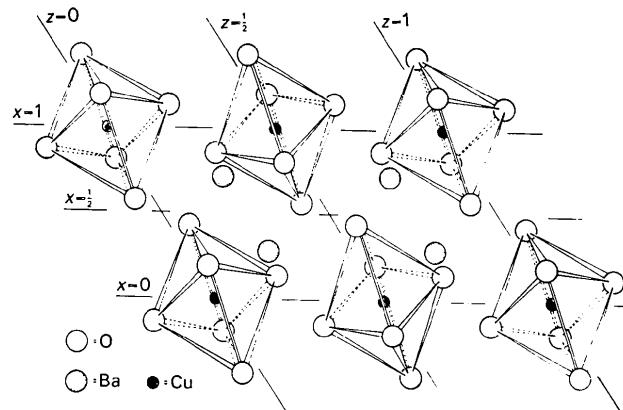


Fig. 1. Arrangement of discrete, distorted  $[\text{Cu}(\text{OH})_6]^{4-}$  octahedra in  $\text{Ba}_2[\text{Cu}(\text{OH})_6]$ , projected along [010].

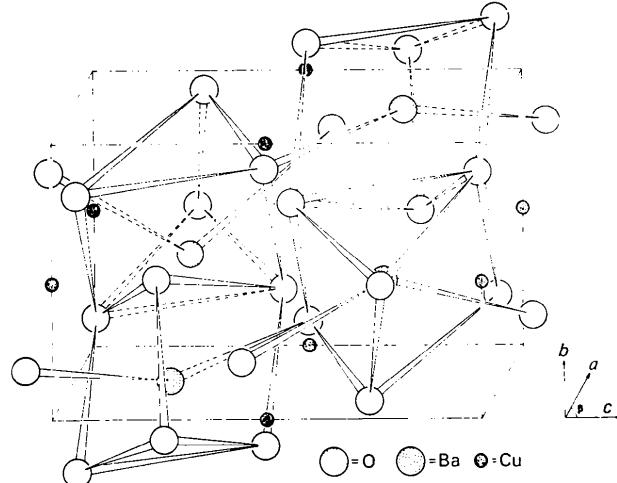


Fig. 2. Three-dimensional view of the  $\text{Ba}_2[\text{Cu}(\text{OH})_6]$  structure, showing the di-capped trigonal prismatic coordination of barium by oxygen.

Table 2. Atomic positional coordinates in  $\text{Ba}_2[\text{Cu}(\text{OH})_6]$

Standard deviations in the last digit calculated in the final least-squares refinement are given in parentheses.

Position	<i>x</i>	<i>y</i>	<i>z</i>	
Ba	4(e)	0.2821 (1)	0.0674 (1)	0.2489 (1)
Cu	2(c)	0.00	0.00	0.50
O(1)	4(e)	0.4327 (8)	0.2575 (6)	0.0586 (4)
O(2)	4(e)	0.1956 (8)	0.7629 (5)	0.0393 (4)
O(3)	4(e)	0.1387 (8)	0.4750 (6)	0.2267 (4)

Table 3. Mean-square amplitudes of vibration ( $\text{\AA}^2 \times 10^{-4}$ )

The temperature factors were calculated from the expression:

$$\exp[-2\pi^2(U_{11}h^2a^*{}^2 + U_{22}k^2b^*{}^2 + U_{33}l^2c^*{}^2 + 2U_{12}hka^*b^* + 2U_{13}hla^*c^* + 2U_{23}klb^*c^*)].$$

	$U_{11}$	$U_{22}$	$U_{33}$	$U_{12}$	$U_{13}$	$U_{23}$
Ba	138 (1)	123 (1)	126 (1)	3 (1)	69 (1)	7 (1)
Cu	133 (4)	128 (4)	104 (3)	24 (4)	61 (3)	1 (3)
O(1)	197 (19)	181 (19)	148 (16)	-30 (17)	94 (15)	-14 (13)
O(2)	178 (19)	145 (17)	166 (16)	-49 (16)	104 (15)	-36 (13)
O(3)	142 (18)	128 (16)	119 (14)	5 (15)	66 (13)	18 (12)

Extinction factor  $g = 4.93 (5) \times 10^{-3}$  (Larson, 1967).

do not share any edges or corners and appear to be held together primarily by barium ions. The arrangement of the isolated copper polyhedra within the structure is shown in Fig. 1.

The coordination of the barium ion may be described as a polyhedron of eight oxygens forming a distorted di-capped trigonal prism (Fig. 2). These prisms are linked together into chains parallel to [010] by edge-sharing of the trigonal base.

The crystal structure can also be described in terms of close-packed hexagonal layers (011) consisting of

Table 4 (*cont.*)

Table 5. Copper(II)-oxygen bond lengths ( $\text{\AA}$ ) of selected (4+2) or (4+1+1) distorted octahedra in recently published crystal structures

Compound	Reference	Central atom	Coordination number	Average of 4 equatorial bonds ( $\text{\AA}$ )	Average of 2 axial bonds ( $\text{\AA}$ )
$\text{Ba}_2[\text{Cu}(\text{OH})_6]$	This work	Cu	4(+2)	1.965	2.805
$\text{CuO}$	Åsbrink & Norrby (1970)	Cu	4(+2)	1.956	2.784
$\text{Cu}_3\text{AsO}_4(\text{OH})_3$	Ghose, Fehlmann & Sundaralingam (1965)	Cu(1)	4+1+1	1.942	2.917
		Cu(2)	4+1+1	1.974	2.642
		Cu(3)	4+1+1	1.964	2.576
$\text{CuHPO}_3 \cdot 2\text{H}_2\text{O}$	Handlovič (1969)	Cu	4+1+1	1.980	2.695
$\text{CuMoO}_4$	Abrahams, Bernstein & Jamieson (1968)	Cu(1)	4+1+1	1.950	2.416
$\text{CuUO}_4$	Siegel & Hoekstra (1968)	Cu(2)	4+2	1.956	2.313
$\text{Cu}_2\text{P}_4\text{O}_{12}$	Laügt, Guitel, Tordjman & Bassi (1972)	Cu	4+2	1.955	2.590
$\text{Cu}_2\text{SO}_3 \cdot \text{CuSO}_3 \cdot 2\text{H}_2\text{O}$	Kierkegaard & Nyberg (1965)	Cu(1)	4+2	1.948	2.471
$\text{CuWO}_4$	Kihlborg & Gebert (1970)	Cu(2)	4+2	1.954	2.369
$\text{CuSO}_4 \cdot 3\text{H}_2\text{O}$	Zahrobsky & Baur (1968)	Cu	4+2	1.979	2.467
				1.976	2.395
				1.961	2.424

This coordination is commonly referred as to (4+2) and is explained as a consequence of the Jahn-Teller theorem. In  $\text{Ba}_2[\text{Cu}(\text{OH})_6]$ , however, the elongation of the axial bonds is considerable and it could be argued that the coordination is actually square planar. Table 5 contains bond lengths in recently published (4+2) and (4+1+1) distorted copper(II)-oxygen octahedra. It shows that the mean of the four equatorial bond lengths of 1.965  $\text{\AA}$  is in very good agreement with other structure determinations. The axial bonds of 2.805  $\text{\AA}$  are at the upper limit of what usually is called a (4+2) coordination; but since these axial atoms lay very close to the expected octahedral positions, as evidenced by the  $\text{O}_{\text{axial}}-\text{Cu}-\text{O}_{\text{eq}}$  angles of 91.7, 88.3, 93.1 and 86.9° (Fig. 3), it is reasonable to describe the coordination of the cupric ion as tetragonal distorted octahedral.

The configuration of the barium-oxygen polyhedron can be best described as follows: six of the coordinating oxygen atoms form a distorted trigonal prism with the barium at its centre; the remaining two oxygen atoms lie just outside the centres of two of the prism faces. For interatomic distances see Table 6.

barium-oxygen distance of 2.809  $\text{\AA}$  compares well with recently reported values.

#### Hydrogen bonding and infrared spectra

The positions of possible hydrogen bonds were inferred from an examination of the distances between oxygen atoms which do not belong to the same coordination polyhedron of copper or barium. The infrared spectrum of  $\text{Ba}_2[\text{Cu}(\text{OH})_6]$  in the  $3\mu$  region consists of a sharp peak at  $3520 \text{ cm}^{-1}$  and a broader band at

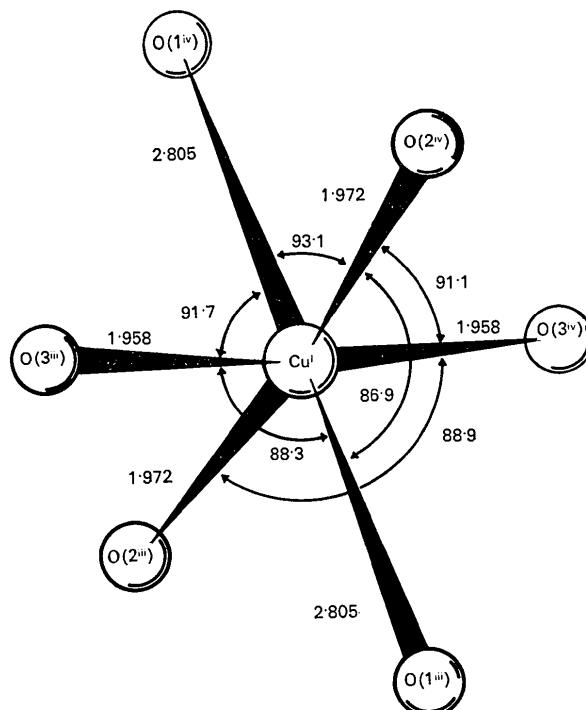


Fig. 3. Interatomic distances ( $\text{\AA}$ ) and bond angles (°) of the  $[\text{Cu}(\text{OH})_6]^{4-}$  polyhedron. The standard deviations  $\sigma$  are 0.004  $\text{\AA}$  for Cu-O bond lengths and 0.15° for O-Cu-O bond angles.

Table 6. Barium-oxygen interatomic distances in  $\text{Ba}_2[\text{Cu}(\text{OH})_6]$  less than 3.20  $\text{\AA}$

$\text{Ba}^{\text{I}}-\text{O}(1^{\text{I}})$	2.826 (5) $\text{\AA}$
$\text{Ba}^{\text{I}}-\text{O}(1^{\text{III}})$	2.642 (3)
$\text{Ba}^{\text{I}}-\text{O}(1^{\text{IV}})$	2.954 (4)
$\text{Ba}^{\text{I}}-\text{O}(2^{\text{I}})$	2.716 (4)
$\text{Ba}^{\text{I}}-\text{O}(2^{\text{II}})$	2.931 (3)
$\text{Ba}^{\text{I}}-\text{O}(2^{\text{III}})$	2.922 (3)
$\text{Ba}^{\text{I}}-\text{O}(3^{\text{I}})$	2.733 (4)
$\text{Ba}^{\text{I}}-\text{O}(3^{\text{II}})$	2.749 (5)
Mean value	2.809

#### Symmetry code

i	$x, y, z;$	iii	$-x, \frac{1}{2}+y, \frac{1}{2}-z$
ii	$-x, -y, -z;$	iv	$x, \frac{1}{2}-y, \frac{1}{2}+z$

As shown in Table 7, this coordination number of eight for barium is rather common, and the mean

Table 7. Mean Ba–O distances as a function of coordination number in recently published crystal structures

Compound	Reference	Central atom	Coordination number	Mean Ba–O distance (Å)
$\text{Ba}_3(\text{VO}_4)_2$	Süss & Buerger (1970)	Ba(1)	6	2.760
$\text{BaTeO}_3 \cdot \text{H}_2\text{O}$	Nielson, Hazell & Rasmussen (1971)	Ba	7	2.743
$\text{BaFe}_2\text{O}_4$	Mitsuda, Mori & Okazaki (1971)	Ba(1)	7	2.807
$\text{BaNaP}_3\text{O}_9$	Martin & Mitschler (1972)	Ba	8	2.801
$\text{Ba}_2[\text{Cu}(\text{OH})_6]$	This work	Ba	8	2.809
$\gamma\text{-Ba}[\text{AlO}(\text{OH})_2]_2$	Glasser & Giovanoli (1972)	Ba	8	2.893
$\alpha\text{-Ba}[\text{AlO}(\text{OH})_2]_2$	Ahmed & Glasser (1970a)	Ba	9	2.912
$\text{Ba}_3(\text{VO}_4)_2$	Süss & Buerger (1970)	Ba(2)	10	2.865
$\text{Ba}_2[\text{Al}_2(\text{OH})_{10}]$	Ahmed & Glasser (1970b)	Ba	11	2.911
$\text{Ba}_6\text{Ti}_{17}\text{O}_{40}$	Tillmanns & Baur (1970)	Ba(1)	11	2.937
$\text{BaFe}_2\text{O}_4$	Mitsuda, Mori & Okazaki (1971)	Ba(2)	11	2.968
$\text{Ba}_6\text{Ti}_{17}\text{O}_{40}$	Tillmanns & Baur (1970)	Ba(3)	12	2.925
$\text{BaC}_2\text{O}_4 \cdot \text{H}_2\text{C}_2\text{O}_4 \cdot 2\text{H}_2\text{O}$	Dusausoy, Protas, Mutin & Watelle (1970)	Ba	12	2.948
$\text{BaTi}_5\text{O}_{11}$	Tillmanns (1969)	Ba	12	2.966

3320 cm<sup>-1</sup>. The sharp peak at 3520 cm<sup>-1</sup> indicates a practically free OH group, while the peak at 3320 cm<sup>-1</sup> shows the presence of an OH group which is influenced by a weak hydrogen bond. There is one reasonable hydrogen-bonding distance of 2.902 Å in the structure between O(1<sup>i</sup>) and O(3<sup>ii</sup>). This distance is in reasonable agreement with the hydrogen bonded O–O length of about 2.85 Å which may be deduced from the infrared peak at 3320 cm<sup>-1</sup> (Schwarzmann 1962). The coordination of O(1) including the hydrogen bonding to O(3) becomes approximately tetrahedral, O(1) being surrounded by three barium atoms and one oxygen atom.

No definite hydrogen bonds seem to exist between O(1)–O(2) and O(2)–O(3).

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